

Soliton Solutions of Nonlinear Fractional Differential Equations via Functional Variable Method

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ABSTRACT

In this paper, the functional variable method is utilized to derive analytical solutions for the $(2 + 1)$ -dimensional time-fractional Zoomeron equation and the space-time fractional modified regularized long-wave equation, based on the Jumarie's modified Riemann-Liouville derivative. The given equations are transformed into nonlinear ordinary differential equations of integer order, which are then solved using the proposed functional variable method, a novel analytical approach. Consequently, several exact solutions are successfully obtained. The results demonstrate that the proposed method is both efficient and easy to implement.

Keywords

$(2 + 1)$ -dimensional time-fractional Zoomeron equation, fractional modified regularized long-wave equation, Jumarie's modified Riemann-Liouville derivative, Functional Variable Method.

1. INTRODUCTION

Fractional differential equations (FDEs) have drawn considerable interest in recent years, as they are increasingly employed to model complex nonlinear phenomena in fields such as physics, biology, mathematics, economics, engineering, and other scientific areas. Many real-world systems are formulated using FDEs, making the study of their exact solutions essential for scientific research. Fractional partial differential equations are a generalized form of nonlinear partial differential equations. They play a crucial role in understanding and analyzing nonlinear phenomena in applied sciences.

In recent decades, numerous researchers have successfully investigated the exact solutions and analytical approximations of nonlinear Fractional differential equations (FDEs). Many semi-analytical and analytical methods, such as modified extended tanh function method [1], modified Kudryashov method [2] and q-homotopy analysis transform method [10], the modified (G'/G) -expansion method [4], the improved Bernoulli sub-equation function method [5], the Kudryashov method [6], the sine-Gordon expansion method [23] [7, 13], the first integral method [9, 10], the extended trial equation method [22], the $(G'/G, 1/G)$ -expansion method [34], the improved F-expansion method [35], the modified simple equation method [15, 16], the improved Bernoulli sub-equation function method [17], the Darboux transformation method [18, 19, 20], fractional subequation method [21], modified trial

equation method [22], etc. have been employed for obtaining new exact solutions of the nonlinear partial differential equations including integer and fractional orders.

A powerful and effective method for finding exact solutions of nonlinear partial differential equations, known as the functional variable method, was proposed by Zerarka et al. in [23, 24]. More recently, Babajanov [25, 26, 27] applied this method to obtain soliton solutions for various differential equations. In [28], Jumarie introduced a modified Riemann-Liouville derivative. Using this type of fractional derivative along with some useful formulas, fractional differential equations can be transformed into integer-order differential equations through variable transformation.

This study aims to highlight the effectiveness of the functional variable method and the modified Riemann-Liouville derivative in solving nonlinear time-fractional differential equations to obtain exact solitary wave solutions, periodic wave solutions, and combined formal solutions.

The rest of this paper is structured as follows. In section 2, we proposed the basic idea of the method for finding exact travelling wave solutions of nonlinear time-fractional differential equations. In section 3, we established the exact travelling wave solution for the $(2+1)$ -dimensional time fractional Zoomeron equation and the space-time fractional modified regularized long-wave equation. Finally, Sections 4 and 5 present the graphical representations of the equations and the conclusions, respectively.

2. JUMARIE'S MODIFIED RIEMANN-LIOUVILLE DERIVATIVE AND THE FUNCTIONAL VARIABLE METHOD

Jumarie's modified Riemann-Liouville derivative of order α is defined as [29]:

$$\mathfrak{D}_t^\alpha f(t) = \begin{cases} \frac{1}{\Gamma(-\alpha)} \int_0^t (t-\zeta)^{-\alpha-1} [f(\zeta) - f(0)] d\zeta, & \alpha < 0, \\ \frac{1}{\Gamma(-\alpha)} \frac{d}{dt} \int_0^t (t-\zeta)^{-\alpha-1} [f(\zeta) - f(0)] d\zeta, & 0 < \alpha < 1, \\ (f^{(n)}(t))^{(\alpha-n)}, & n \leq \alpha \leq n+1, n \geq 1 \end{cases}$$

where $f : R \rightarrow R, x \rightarrow f(x)$ denote a continuous (but not necessarily differentiable) function. We list some important properties for Jumarie's fractional derivative as

$$\mathfrak{D}_t^\alpha t^r = \frac{\Gamma(1+r)}{\Gamma(1+r-\alpha)} t^{r-\alpha}, r > 0,$$

$$\mathfrak{D}_t^\alpha c = 0, \quad c = \text{constant},$$

$$\mathfrak{D}_t^\alpha [af(t) + bg(t)] = a\mathfrak{D}_t^\alpha f(t) + b\mathfrak{D}_t^\alpha g(t),$$

$$\mathfrak{D}_t^\alpha (f(t)g(t)) = g(t)\mathfrak{D}_t^\alpha f(t) + f(t)\mathfrak{D}_t^\alpha g(t),$$

$$\mathfrak{D}_t^\alpha f(g(t)) = f'_g(g(t))\mathfrak{D}_t^\alpha g(t) = \mathfrak{D}_g^\alpha f(g(t))(g'(t))^\alpha.$$

Drawing inspiration from the works of Lu [30] and Zerarka et al. [31], we present the functional variable method for obtaining exact solutions of nonlinear time-fractional differential equations, as described below.

Let us consider the time-fractional differential equation with independent variables t, x, y, z, \dots and a dependent variable u

$$\mathcal{P}(u, \mathfrak{D}_t^\alpha u, u_x, u_y, u_z, \mathfrak{D}_t^{2\alpha} u, u_{xy}, u_{yz}, u_{xz}, \dots) = 0 \quad (1)$$

where the subscript denotes partial derivative. Using the variable transformation

$$u(t, x, y, z, \dots) = U(\zeta), \quad \zeta = a_1 x + a_2 y + a_3 z + \dots \pm \frac{\omega t^\alpha}{\Gamma(1+\alpha)}$$

where a_i and ω are constants to be determined later; the fractional differential equation (1) is reduced to an ordinary differential equation (ODE)

$$\mathcal{Q}(U, U_\zeta, U_{\zeta\zeta}, U_{\zeta\zeta\zeta}, \dots) = 0 \quad (2)$$

Then we make a transformation in which the unknown function U is considered as a functional variable in the form

$$U_\zeta = F(U) \quad (3)$$

and some successive derivatives of U are

$$\begin{aligned} U_{\zeta\zeta} &= \frac{1}{2}(F^2(U))' \\ U_{\zeta\zeta\zeta} &= \frac{1}{2}(F^2(U))'' \sqrt{F^2(U)} \\ U_{\zeta\zeta\zeta\zeta} &= \frac{1}{2} \left[(F^2(U))''' F^2 + \frac{1}{2}(F^2(U))'' (F^2(U))' \right] \end{aligned} \quad (4)$$

and so on, where $' = d/dU$

Substituting (4) into (2), we reduce the ODE (2) in terms of U, F and its derivatives as

$$\mathcal{R}(U, F, F', F'', F''', \dots) = 0 \quad (5)$$

Equation (5) is particularly important because it admits analytical solutions for a large class of nonlinear wave-type equations. After integration, eq. (5) provides the expression for F , and this together with eq. (3) give relevant solutions to the original problem. In order to illustrate how the method works, we examine some examples treated by other approaches. This is discussed in the following section.

3. APPLICATIONS

3.1 The (2+1)-dimensional time fractional Zoomeron equation

In this subsection, we employ the Functional Variable Method(FVM) to obtain general exact solutions of the time fractional Zoomeron equation, which has been proposed as follows

$$\mathfrak{D}_{tt}^{2\alpha} \left(\frac{u_{xy}}{u} \right) - \left(\frac{u_{xy}}{u} \right)_{xx} + 2\mathfrak{D}_t^\alpha (u^2)_x = 0, \quad t > 0, 0 < \alpha \leq 1 \quad (6)$$

where $u = u(x, y, t)$ is the amplitude of the relevant wave mode. Equation 6 plays a significant role in describing the evolution of a single scalar field and serves as a useful model for showcasing novel phenomena associated with boomerons and trappons [32]. This equation was initially introduced by Calogero and Degasperis [33]. In the past, many researchers have studied both the integer and fractional forms of the Zoomeron equation [34, 35, 36, 37, 38, 39].

Let us consider the travelling wave solutions of Eq. (6), and we perform the transformation:

$$u(x, y, t) = U(\zeta), \quad \zeta = ax + by - \frac{\omega t^\alpha}{\Gamma(1+\alpha)} \quad (7)$$

where a, b and ω are constants.

Substituting transformation (7) with (1) into Eq.(6). Then Eq. (6) can be reduced to an ODE in the following form:

$$ab\omega^2 \left(\frac{U''}{U} \right)'' - a^3 b \left(\frac{U''}{U} \right)' - 2a\omega (U^2)'' = 0 \quad (8)$$

Integrating Eq. (8) twice with respect to ζ , we get

$$ab(\omega^2 - a^2)U'' - 2a\omega U^3 - kU = 0 \quad (9)$$

where, primes denote differentiation with respect to ζ and k is a non zero constant of integration, while the second constant of integration is vanishing.

Substituting eq. 4 into eq. 9, we get

$$\frac{ab(\omega^2 - a^2)}{2} (F^2(U))' = kU + 2a\omega U^3 \quad (10)$$

Integrating the Eq. 10 with respect to U with zero constants of integration, we have

$$F(U) = \pm \sqrt{\frac{\omega}{b(\omega^2 - a^2)}} U \sqrt{\frac{k}{a\omega} + U^2} \quad (11)$$

From eq. 4 and 11 we deduce that

$$\int \frac{dU}{U \sqrt{\frac{k}{a\omega} + U^2}} = \pm \sqrt{\frac{\omega}{b(\omega^2 - a^2)}} (\zeta + \zeta_0) \quad (12)$$

where ζ_0 is a constant of integration. After integrating 12, we have the following exact solutions, for $\frac{\omega}{b(\omega^2 - a^2)} > 0$:

$$u_{1,1}(x, t) = \pm \sqrt{\frac{k}{a\omega}} \operatorname{csch} \left[\sqrt{\frac{k}{ab(\omega^2 - a^2)}} \left(ax + by - \frac{\omega t^\alpha}{\Gamma(1+\alpha)} + \zeta_0 \right) \right] \quad (13)$$

$$u_{1,2}(x, t) = \pm \sqrt{\frac{-k}{a\omega}} \operatorname{sech} \left[\sqrt{\frac{k}{ab(a^2 - \omega^2)}} \left(ax + by - \frac{\omega t^\alpha}{\Gamma(1 + \alpha)} + \zeta_0 \right) \right] \quad (14)$$

For $\frac{\omega}{b(\omega^2 - a^2)} < 0$, we obtain periodic solutions as follows:

$$u_{1,3}(x, t) = \pm \sqrt{\frac{k}{a\omega}} \csc \left[\sqrt{\frac{k}{ab(a^2 - \omega^2)}} \left(ax + by - \frac{\omega t^\alpha}{\Gamma(1 + \alpha)} + \zeta_0 \right) \right] \quad (15)$$

$$u_{1,4}(x, t) = \pm \sqrt{\frac{-k}{a\omega}} \sec \left[\sqrt{\frac{k}{ab(a^2 - \omega^2)}} \left(ax + by - \frac{\omega t^\alpha}{\Gamma(1 + \alpha)} + \zeta_0 \right) \right] \quad (16)$$

3.2 The space-time fractional modified regularized long-wave equation

We next consider the space-time fractional modified regularized long-wave equation (mRLW) [40]

$$\mathfrak{D}_t^\alpha u + a\mathfrak{D}_x^\alpha u + bu^2\mathfrak{D}_x^\alpha u - c\mathfrak{D}_t^\alpha \mathfrak{D}_x^{2\alpha} u = 0, \quad 0 < \alpha \leq 1 \quad (17)$$

where a, b, c are arbitrary constants. The regularized long wave model is a fundamental and significant model in nonlinear science, first introduced by Peregrine [41]. This model proves to be valuable in explaining various nonlinear phenomena across different scientific and engineering fields, including pressure waves in liquids, gas bubbles, ion-acoustic and hydrodynamic waves in plasma, phonon packets in nonlinear crystals, and longitudinal dispersive waves in elastic rods, among others.

Let us introduce the following transformations:

$$u(x, t) = U(\zeta), \quad \zeta = \frac{x^\alpha}{\Gamma(1 + \alpha)} + \omega \frac{t^\alpha}{\Gamma(1 + \alpha)} \quad (18)$$

Here ω is a non-zero constant. Substituting 18 into 17, we get the following ODE:

$$(\omega + a)U' + bU^2U' - c\omega U''' = 0 \quad (19)$$

Further, integrating 19 with respect to ζ , we get

$$(\omega + a)U + \frac{b}{3}U^3 - c\omega U'' = 0 \quad (20)$$

Substituting eq. 4 into eq. 20, we get

$$\frac{c\omega}{2}(F^2(U))' = (\omega + a)U + \frac{b}{3}U^3 \quad (21)$$

Integrating the Eq. 21 with respect to U with zero constants of integration, we have

$$F(U) = \pm \sqrt{\frac{b}{6c\omega}} U \sqrt{\frac{6(\omega + a)}{b} + U^2} \quad (22)$$

From eq. 4 and 22 we deduce that

$$\int \frac{dU}{U \sqrt{\frac{6(\omega + a)}{b} + U^2}} = \pm \sqrt{\frac{b}{6c\omega}} (\zeta + \zeta_0) \quad (23)$$

where ζ_0 is a constant of integration. After integrating 23, we have the following exact solutions, for $\frac{b}{c\omega} > 0$:

$$u_{2,1}(x, t) = \pm \sqrt{\frac{6(\omega + a)}{b}} \operatorname{csch} \left[\sqrt{\frac{(\omega + a)}{c\omega}} \left(\frac{x^\alpha}{\Gamma(1 + \alpha)} + \omega \frac{t^\alpha}{\Gamma(1 + \alpha)} + \zeta_0 \right) \right] \quad (24)$$

$$u_{2,2}(x, t) = \pm \sqrt{\frac{-6(\omega + a)}{b}} \operatorname{sech} \left[\sqrt{\frac{(\omega + a)}{c\omega}} \left(\frac{x^\alpha}{\Gamma(1 + \alpha)} + \omega \frac{t^\alpha}{\Gamma(1 + \alpha)} + \zeta_0 \right) \right] \quad (25)$$

For $\frac{b}{c\omega} < 0$, we obtain periodic solutions as follows:

$$u_{2,3}(x, t) = \pm \sqrt{\frac{6(\omega + a)}{b}} \csc \left[\sqrt{\frac{-(\omega + a)}{c\omega}} \left(\frac{x^\alpha}{\Gamma(1 + \alpha)} + \omega \frac{t^\alpha}{\Gamma(1 + \alpha)} + \zeta_0 \right) \right] \quad (26)$$

$$u_{2,4}(x, t) = \pm \sqrt{\frac{-6(\omega + a)}{b}} \sec \left[\sqrt{\frac{-(\omega + a)}{c\omega}} \left(\frac{x^\alpha}{\Gamma(1 + \alpha)} + \omega \frac{t^\alpha}{\Gamma(1 + \alpha)} + \zeta_0 \right) \right] \quad (27)$$

4. GRAPHICAL REPRESENTATION

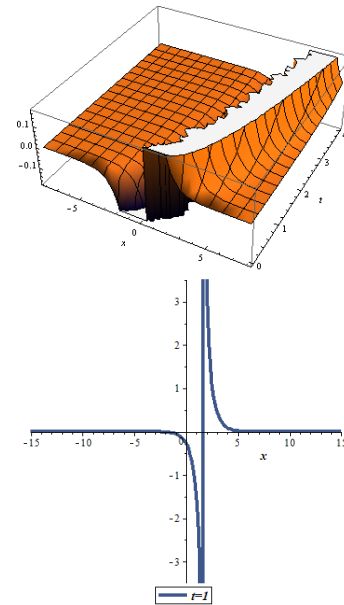
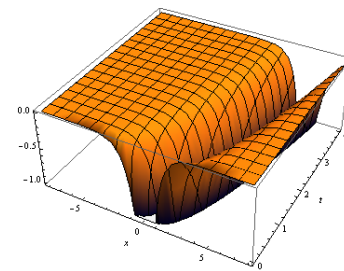


Fig.1: 3D and 2D plot of $u_{1,1}(x, t)$ given in Eq. 13 with $a = 1, b = 1, \omega = 2, \alpha = 0.5, k = -8, y = 0, \zeta_0 = 0$.



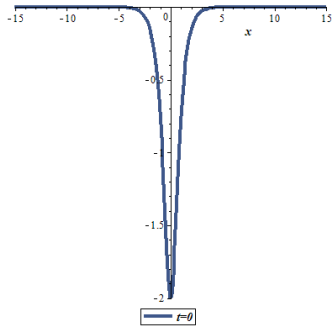


Fig.2: 3D and 2D plot of $u_{1,2}(x, t)$ given in Eq. 14 with $a = 1, b = 1, \omega = 2, \alpha = 0.5, k = -8, y = 0, \zeta_0 = 0$.

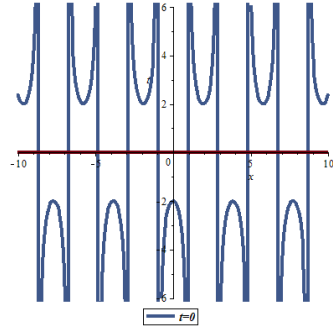


Fig.4: 3D and 2D plot of $u_{1,4}(x, t)$ given in Eq. 16 with $a = 1, b = 1, \omega = 2, \alpha = 0.5, k = -8, y = 0, \zeta_0 = 0$.

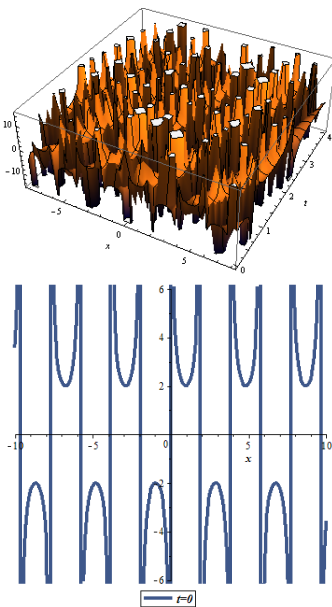


Fig.3: 3D and 2D plot of $u_{1,3}(x, t)$ given in Eq. 15 with $a = 1, b = 1, \omega = 2, \alpha = 0.5, k = -8, y = 0, \zeta_0 = 0$.

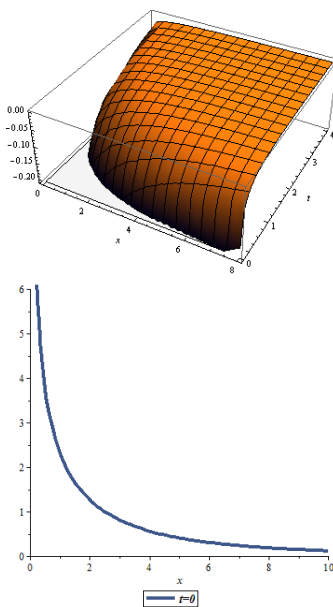
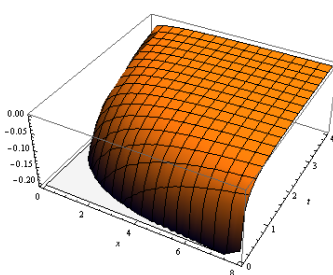
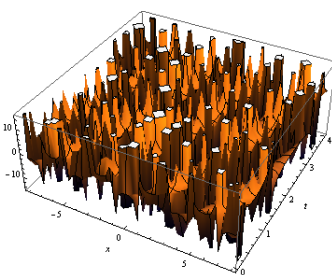


Fig.5: 3D and 2D plot of $u_{2,1}(x, t)$ given in Eq. 24 with $a = b = c = 1, \omega = 2, \alpha = 0.5, \zeta_0 = 0$.



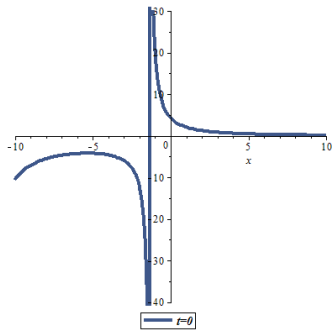


Fig.6: 3D and 2D plot of $u_{2,2}(x, t)$ given in Eq. 25 with $a = b = c = 1, \omega = 2, \alpha = 0.5, \zeta_0 = 0$.

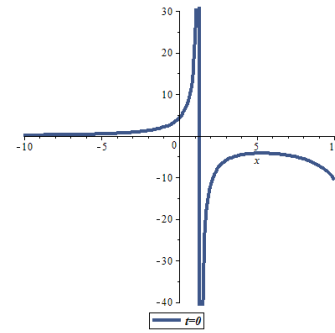


Fig.8: 3D and 2D plot of $u_{2,4}(x, t)$ given in Eq. 27 with $a = b = c = 1, \omega = 2, \alpha = 0.5, \zeta_0 = 0$.

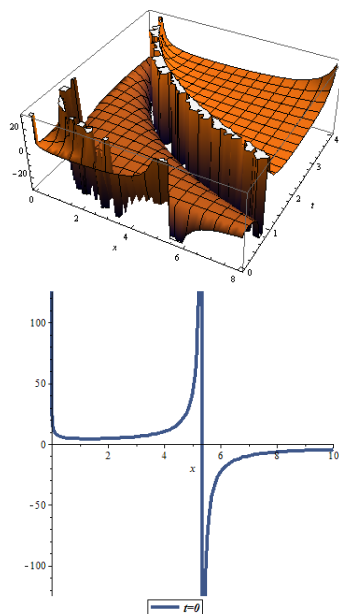
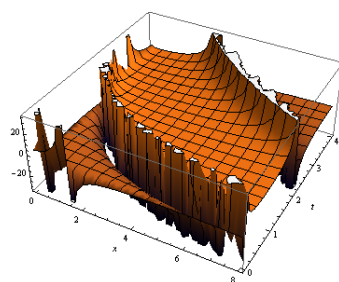


Fig.7: 3D and 2D plot of $u_{2,3}(x, t)$ given in Eq. 26 with $a = b = c = 1, \omega = 2, \alpha = 0.5, \zeta_0 = 0$.



We have provided graphs of solitary waves for equations 6 and 17, created by selecting suitable values for the relevant parameters, to better understand the mechanisms behind the original physical phenomena. Graphical representations are a powerful tool for communication, effectively illustrating the solutions to these problems. These graphs are displayed in Fig.1 to Fig.8. Solitary and periodic wave solutions are significant types of solutions for nonlinear partial differential equations, as many such equations exhibit various solitary wave solutions. Solitons, a specific class of solutions to nonlinear partial differential equations with weak linearity, are frequently used to model physical systems. The existence of periodic traveling waves generally depends on the parameter values in the mathematical equations, with these parameters affecting both amplitude and velocity. A soliton is a self-sustaining wave packet that retains its shape while traveling at a constant velocity.

5. CONCLUSION

In this paper, the functional variable method and the modified Riemann-Liouville derivative are introduced for solving the $(2 + 1)$ -dimensional time-fractional Zoomeron equation and the space-time fractional modified regularized long-wave equation. It is anticipated that the solutions obtained in this study will be valuable for further exploration of complex nonlinear physical phenomena. This method provides a promising approach for solving a wide range of fractional partial differential equations and serves as a reliable technique for handling nonlinear fractional differential equations. Additionally, this method is straightforward, concise, and especially well-suited for computer implementation. The algebraic complexities and extensive calculations were effectively managed using the symbolic computation software Mathematica. As a result, this approach can be extended to tackle nonlinear problems in soliton theory and other related areas.

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